
New Physics Searches at the Tevatron in 2015

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**It is hard to compete with the LHC!
(and even harder with LHC+ILC...)**

Still, the Tevatron has smaller backgrounds and more \bar{q} per collision.

→ there are some examples of physics beyond the standard model where the Tevatron would be useful after the LHC.

Warning: It is hard to say at this point how likely these scenarios are.

Light Stop in the MSSM

Demina, Lykken, Matchev, Nomerotski, hep-ph/9910275

Typically: $\tilde{t} \rightarrow t\chi_1^0$

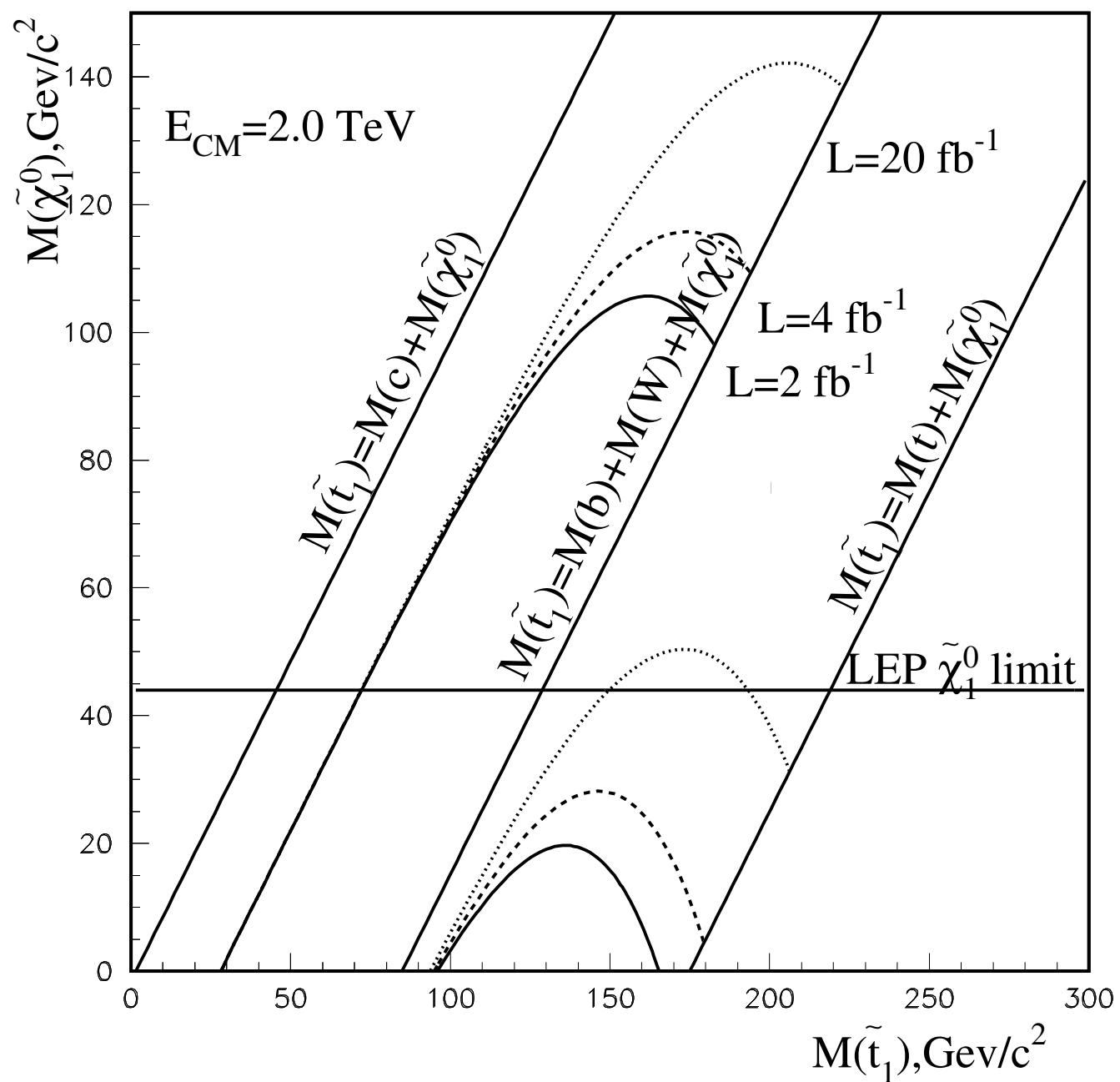
For $m_{\tilde{t}} < m_t + m_\chi$:

$\tilde{t} \rightarrow c\chi_1^0$ is the dominant decay mode

Stop searches in Run II: $\tilde{t}\tilde{t}^\star \rightarrow c\bar{c} E_T$

Challenging at the LHC because of large backgrounds.

$\tilde{t}_1 \rightarrow c \tilde{\chi}_1^0$ or $\tilde{t}_1 \rightarrow b W \tilde{\chi}_1^0$



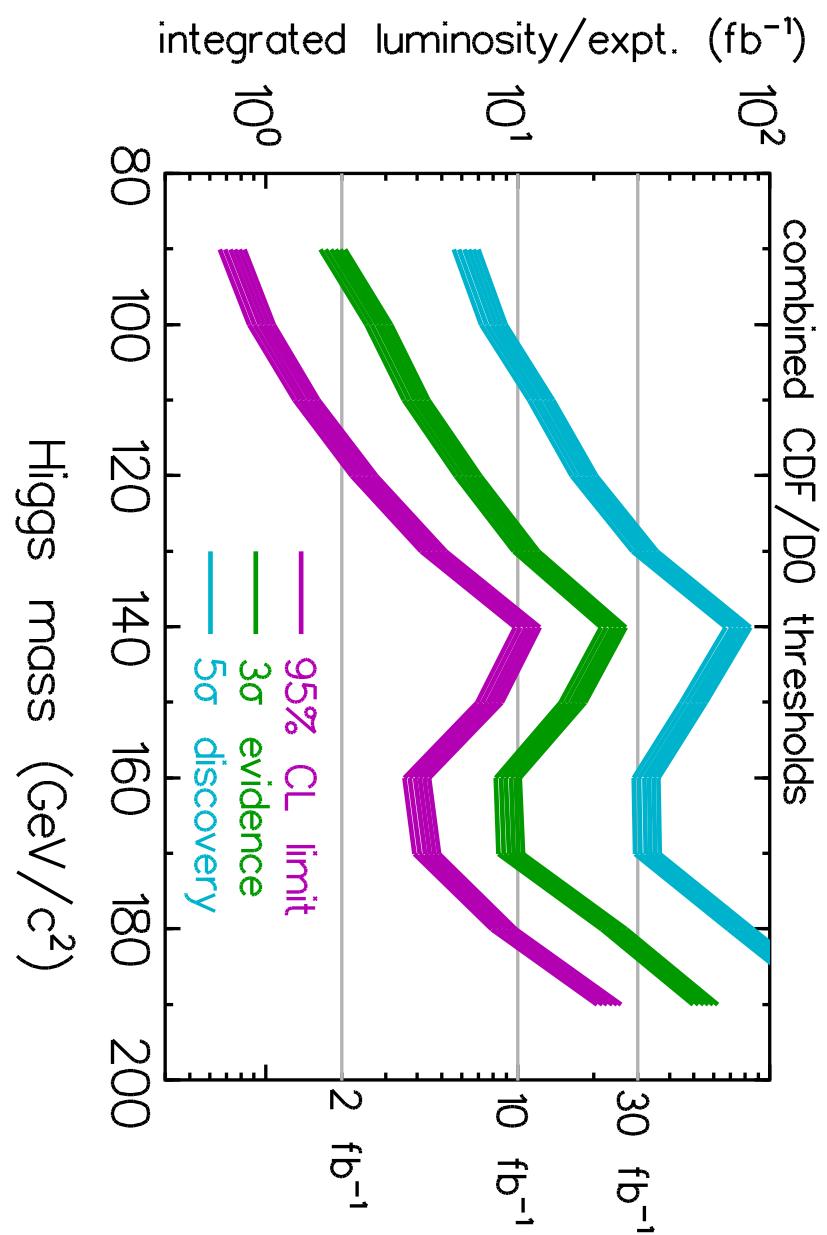
Higgs bosons

For $115 \text{ GeV} < M_h < 130 \text{ GeV}$:

the Tevatron may measure the $h\bar{b}b$ coupling

Associated production ($q\bar{q} \rightarrow W^* \rightarrow hW$) with $h \rightarrow \bar{b}b$

At the LHC: huge $\bar{b}b$ background.



Z' bosons

M. Carena, A. Daleo, B. Dobrescu, T. Tait: hep-ph/0408098

New $U(1)_z$ gauge symmetry

Assume:

- generation-independent charges
- quark masses from standard model Yukawa couplings
- no gauge anomalies

Fermion gauge charges:

	$SU(3)$	$SU(2)$	$U(1)_Y$	$U(1)_{B-xL}$	$U(1)_{q+xu}$	$U(1)_{10+x\bar{5}}$	$U(1)_{d-xu}$
q_L	3	2	1/3	1/3	1/3	1/3	0
u_R	3	1	4/3	1/3	$x/3$	-1/3	$-x/3$
d_R	3	1	-2/3	1/3	$(2-x)/3$	$-x/3$	1/3
l_L	1	2	-1	$-x$	-1	$x/3$	$(-1+x)/3$
e_R	1	1	-2	$-x$	$-(2+x)/3$	-1/3	$x/3$
ν_R	1	1	0	-1	$(-4+x)/3$	$(-2+x)/3$	$-x/3$
ν'_R	$-1-x/3$.
ψ_L^l	1	2	-1	-1	.	$-(1+x)/3$	$-2x/5$
ψ_R^l	.	.	.	$-x$.	2/3	$(-1+x/5)/3$
ψ_L^e	1	1	-2	-1	.	.	.
ψ_R^e	.	.	.	$-x$.	.	.
ψ_L^d	3	1	-2/3	.	.	-2/3	$(1-4x/5)/3$
ψ_R^d	$(1+x)/3$	$x/15$

In general: Z' mixes with Z

\Rightarrow LEP I requires $M_{Z'} \gtrsim 2$ TeV

Special case: $SU(3)_C \times SU(2)_W \times U(1)_Y \times \textcolor{red}{U(1)_{B-L}}$

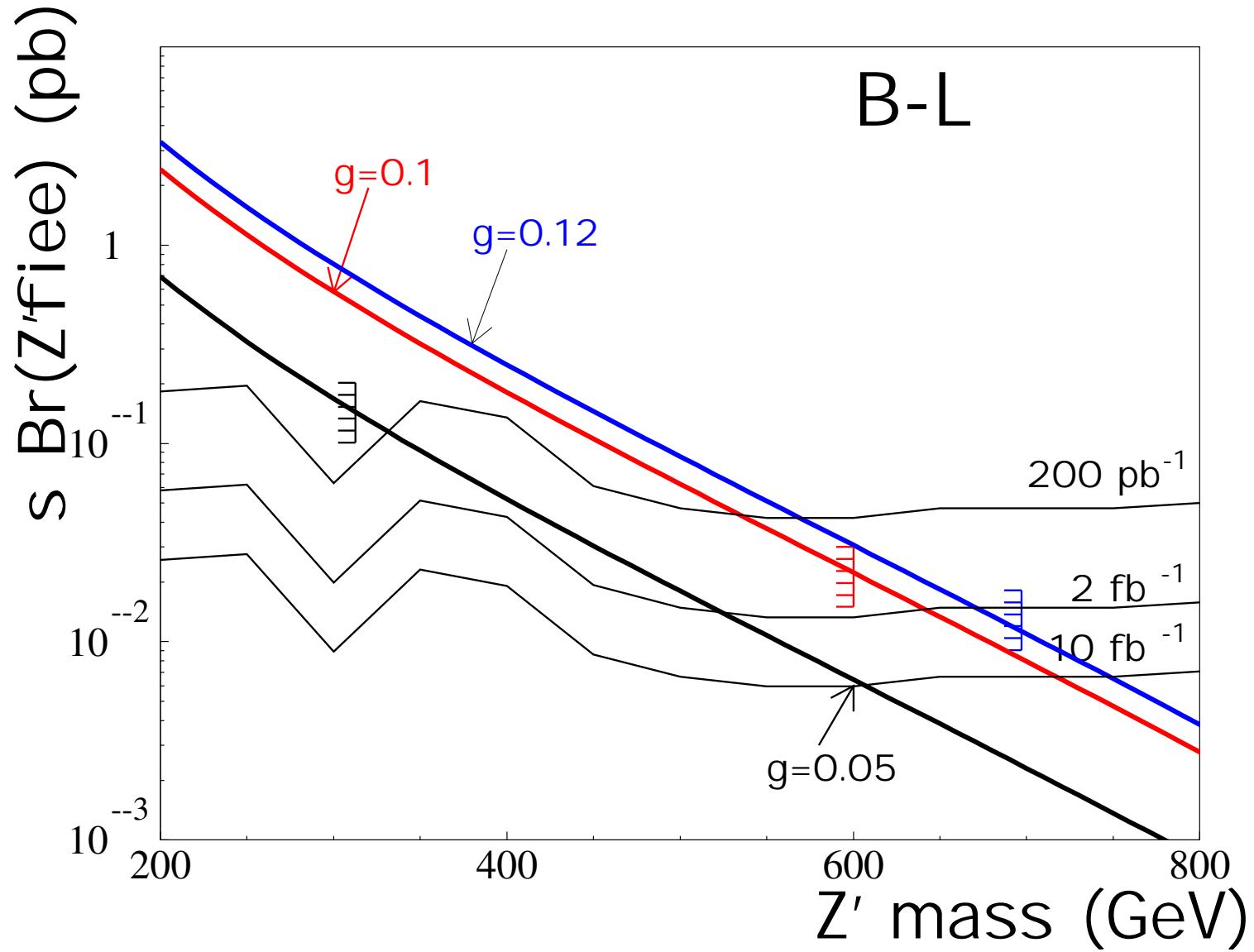
$$z_q = z_u = z_d = -\frac{z_l}{3} = -\frac{z_e}{3} = -\frac{z_\nu}{3} \implies z_H = 0$$

No Z_{B-L} - Z mixing at tree level!

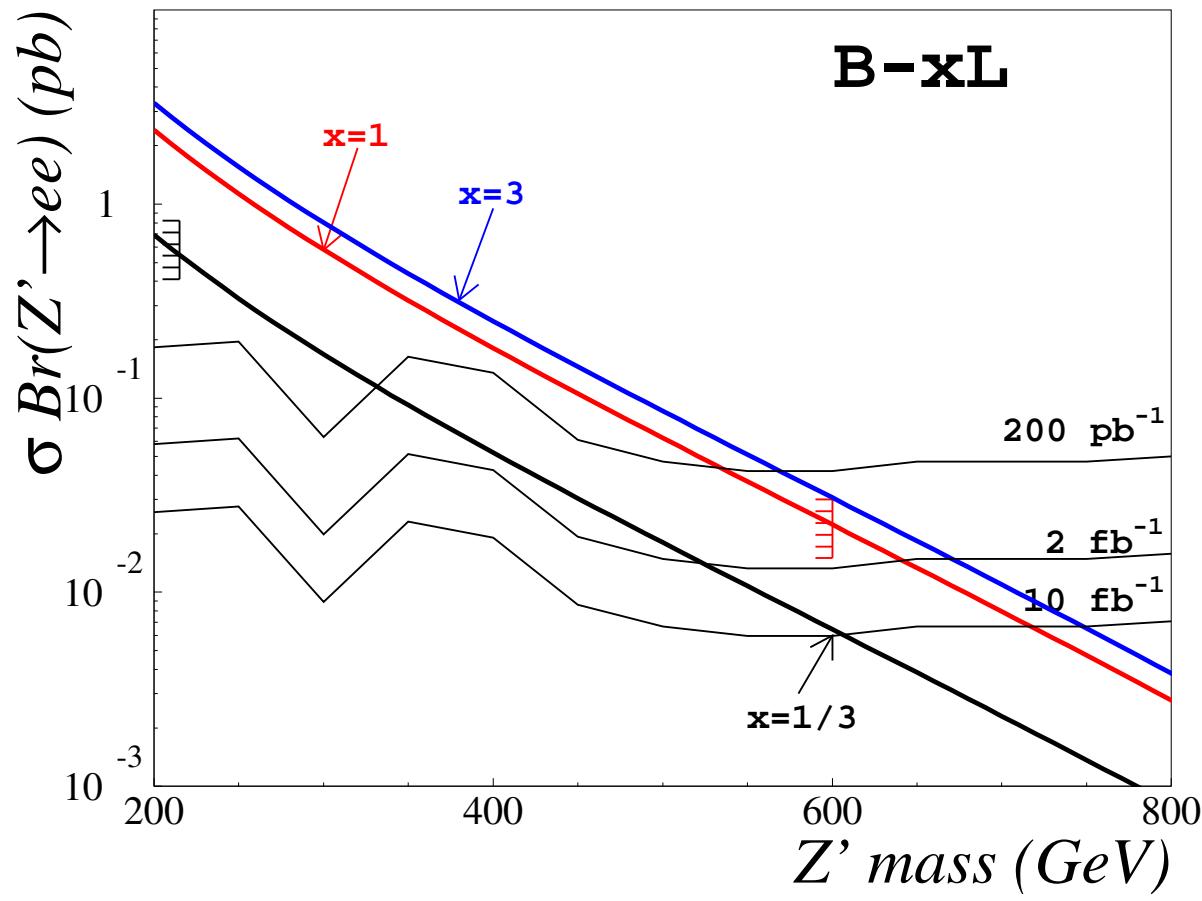
*Best bounds on $z_l g_z$ come from limits on direct production
at the Tevatron and at LEP II*

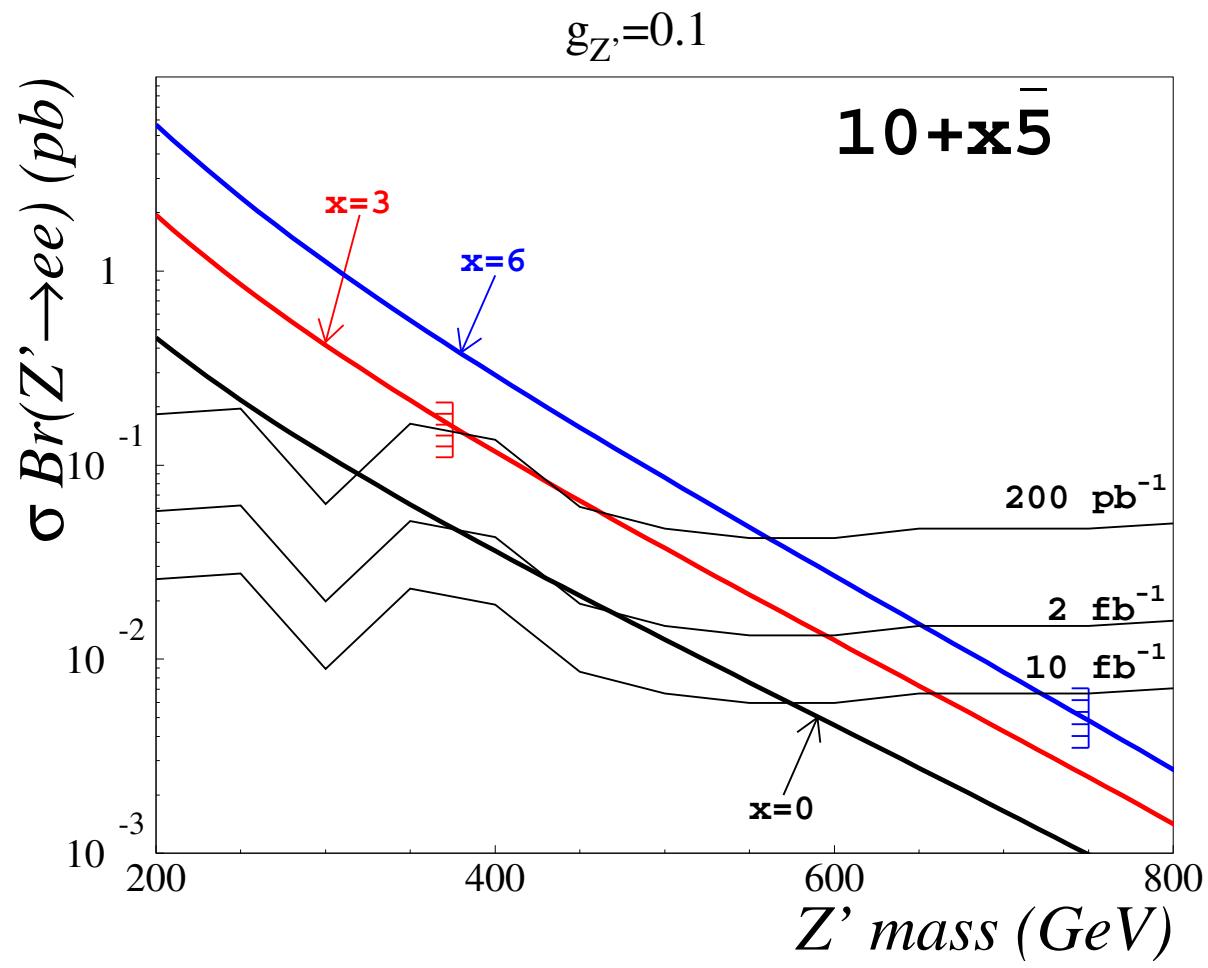
$$\sigma^{Z'} = \frac{\pi}{48 s} W_{Z'} \left(\frac{M_{Z'}^2}{s}, M_{Z'}^2 \right) \mathbf{Br}(Z' \rightarrow l^+ l^-),$$

Z' searches at the Tevatron:



$g_{Z'}=0.1$



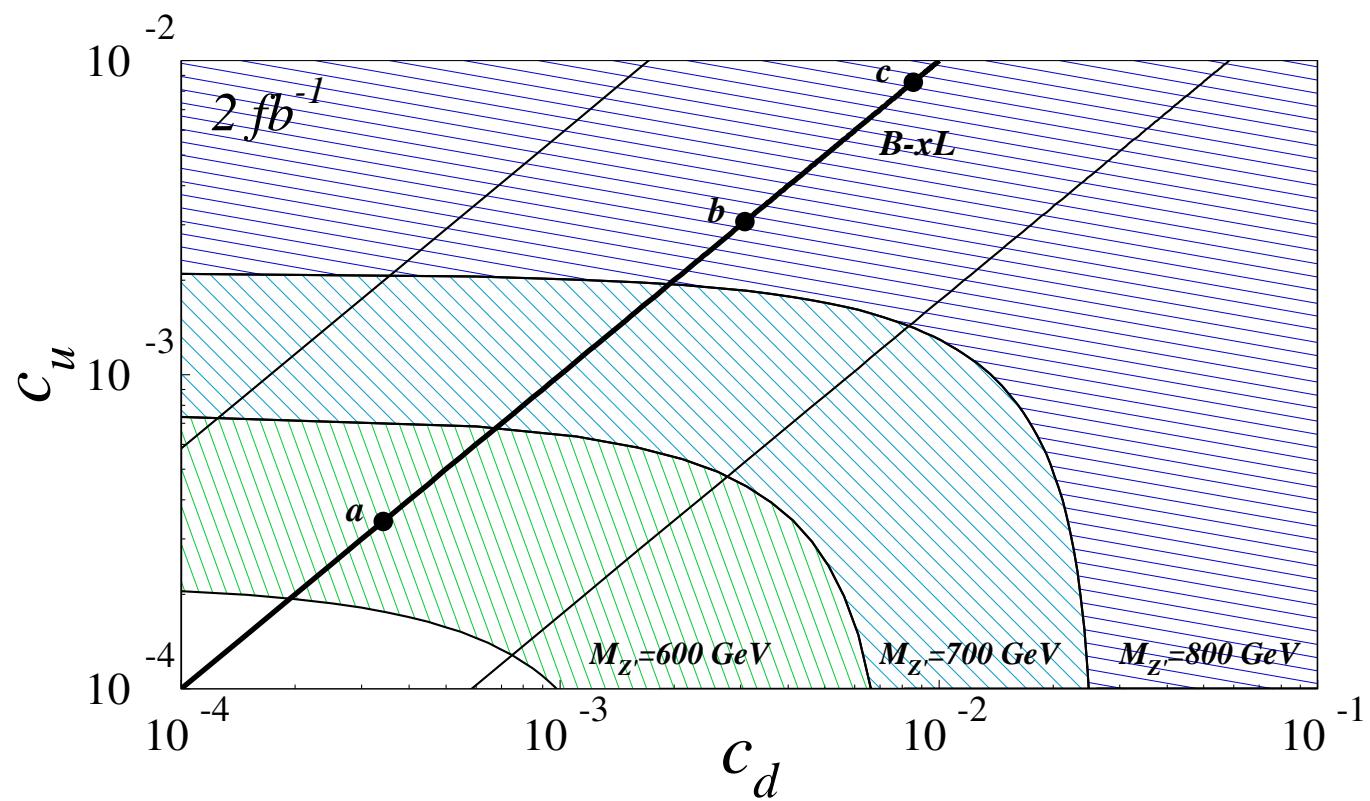


A user-friendly parametrization:

$$\sigma^{Z'} = \frac{\pi}{48 s} \left[c_u w_u \left(\frac{M_{Z'}^2}{s}, M_{Z'}^2 \right) + c_d w_d \left(\frac{M_{Z'}^2}{s}, M_{Z'}^2 \right) \right]$$

$$w_{u(d)} = \sum_q \int_0^1 dx_1 dx_2 dz \{ f_{q/P}(x_1, M_{Z'}^2) f_{\bar{q}/\bar{P}}(x_2, M_{Z'}^2) \Delta_{qq}(z, M_{Z'}^2) \\ + f_{g/P}(x_1, M_{Z'}^2) (f_{q/\bar{P}}(x_2, M_{Z'}^2) + f_{\bar{q}/\bar{P}}(x_2, M_{Z'}^2)) \Delta_{gq}(z, M_{Z'}^2) \\ + (x_1 \leftrightarrow x_2, P \leftrightarrow \bar{P}) \} \delta \left(\frac{M_{Z'}^2}{s} - x x_1 x_2 \right)$$

$$c_{u,d} = g_z^2 (z_q^2 + z_{u,d}^2) \mathbf{Br}(Z' \rightarrow l^+ l^-)$$



Leptophobic Z' , e.g., $U(1)_B$

- No significant bound from LEP or ILC
- LHC has large dijet background
- The Tevatron with high luminosity has probably the best chance of discovering a leptophobic Z' .

Conclusion

After the ATLAS and CMS will start presenting physics results it will be easier to guess what a high luminosity Tevatron will be able to do in 2015!

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